Final state radiation (FSR) impacts high and ultrahigh energy neutrino observations

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Based on arXiv: 2403.07984 Ryan Plestid (Caltech), Bei Zhou (Fermilab & KICP)

More than half a century after the establishment of the quantum electrodynamics, it still has a radiative correction of as large as 25% to be studied.

And it has also been overlooked by current experiments on HE and UHE neutrinos.

Why do we study HE&UHE neutrinos

- Astrophysics (highlighted by astro2020): Origin of HE/UHE astrophysical neutrinos
 - Sources of HE/UHE cosmic rays (> 60-year problem)
 - Cosmic particle acceleration, propagation
 - Cosmic gamma ray sources, hadronic vs leptonic mechanism
 - Dense astrophysical environments
 - Essential for multi-messenger astrophysics



- Particle physics (highlighted by P5 report):
 - Neutrino interactions in the SM (Deep-inelastic scattering, W-boson production, Glashow resonance, final state radiation, etc.)
 - Measure neutrino mixing parameters
 - Test BSM (ν portal to DM, new ν interactions, sterile ν, magnetic moment, etc.)

Lots of HE/UHE nu telescopes running or to build

HE neutrino telescopes (~100 GeV--100 PeV)

Detector	Size	Status
IceCube	1 km³	Running for ~14 yrs
KM3NET	1 km³	Running, constructing
Baikal-GVD	1 km³	Running, constructing
P-ONE	multi-km³	Proposed
IceCube-Gen2	7.9 km³	Proposed
TRIDENT	7.5 km³	Prototype
Etc		

Laboratory HE nu experiments (~10 GeV--5 TeV)

Detector	Size	Status
FASERv	Neutrino beam	Running
SND@LHC	Neutrino beam	Running
FASERv2	Neutrino beam	Proposed
AdvSND@LHC	Neutrino beam	Proposed
FLArE	Neutrino beam	Proposed

UHE neutrino telescopes (>~100 PeV)

Size	Status
	Finished
	Running
	Running
	Constructing
	Constructing
	Prototype
	Prototype
	Proposed
	Prototype
	Size

2203.08096, Ackermann,, <u>BZ</u> (Snowmass) for a complete list

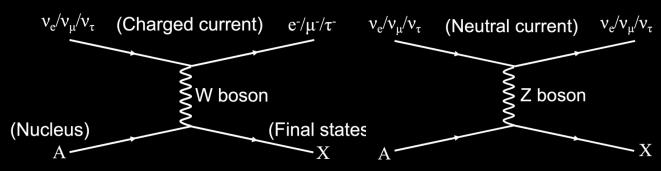
Increasing statistics requires studies of HE/UHE nu interactions

- Neutrino interactions are the cornerstone of all kinds of neutrino-related measurements
 - Astrophysics: energy spectrum, flavor composition, arrival direction, etc.
 - Particle physics: mixing parameters; all BSM studies contingent on well-understood SM interactions
- Help us to find new event classes: useful for both astrophysics and particle physics studies
 - E.g., dimuons for high-energy neutrino detection (2110.02974 <u>BZ</u>, Beacom).
- Neutrino(-nucleus) interaction theory is interesting (and sometimes difficult):
 - Neutrino only has weak interactions, but neutrino interaction studies involves much more
 - Weak, electroweak
 - QED (e.g., final state ration, W-boson and trident production)
 - Strong interactions: QCD (parton distribution functions), nuclear model, resonance prod., etc.
 - (Also detection physics because you need to detect them.)

Overview of HE&UHE neutrino interactions

Deep inelastic scattering (DIS) dominates

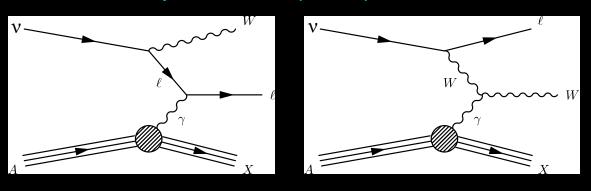
(as good as ~1% precision)



Gandhi+ 96&97, Connolly+ 11, Cooper-Sarkar+ 11, Bertone+ 16, etc.

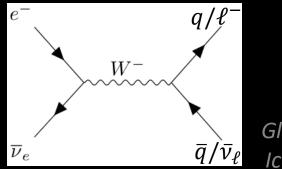
Most recent: Xie, et al. 2303.13607

W-boson production (WBP) is subdominant



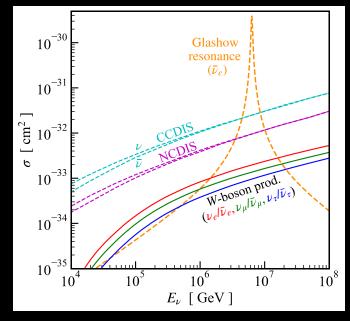
(Seckel 1997, Alikhanov 2015, <u>BZ</u>, Beacom, 1910.08090)

Glashow resonance for $\bar{\nu}_e$



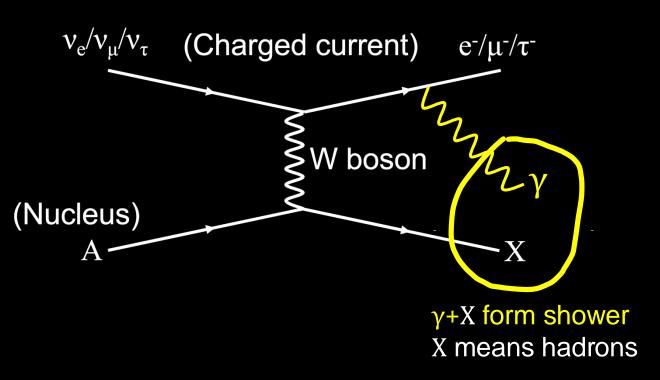
Glashow 1960 IceCube 2021

Cross sections



(BZ, Beacom, 1910.10720)

Our work: final state radiation (FSR) on top of neutrino CC DIS

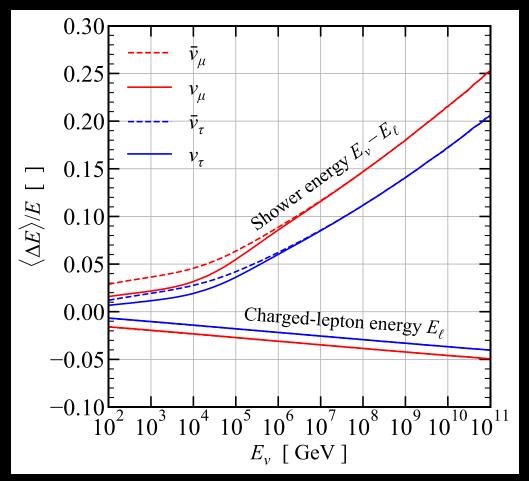


Effect on total xsec: small (~1%, c.f.).

Effects on the differential xsec: big, due to the kinematic logs.

→ So, it affects observation if charged lepton and shower are separate.

FSR impacts the energies of the final states from HE/UHE interactions



Correction increases with energy, up to 25%(!)

Correction on $v\mu > v\tau$, cuz $m_\mu < m_\tau$

Correction on shower > charged lepton

Correction on shower further enhanced by 10—20% due to light yields from EM shower > hadronic shower

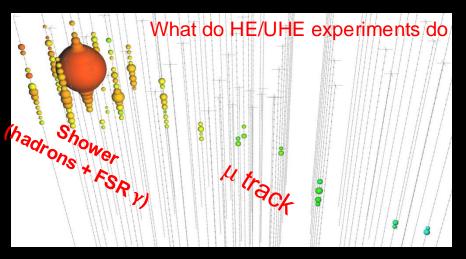
Difference between nu and nubar

Photon takes energy from the charged lepton to the shower

(Plestid, <u>BZ</u>, 2303.08984)

Final state radiation impacts the inelasticity (y) measurements

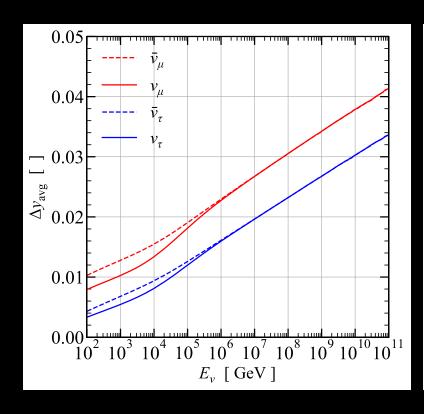
$$y_{\rm QCD} \equiv \frac{E_X}{E_\nu} = \frac{E_\nu - E_\ell}{E_\nu}$$

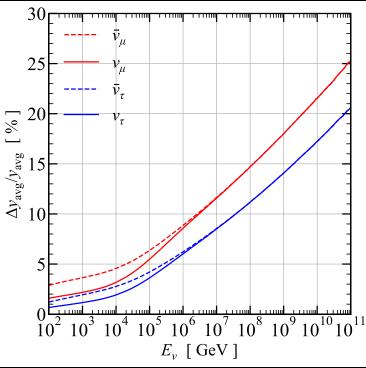


$$y_{\rm exp} \equiv \frac{E_{\rm shower}}{E_{\rm track} + E_{\rm shower}} = y_{\rm QCD} + \frac{E_{\gamma}}{E_{\nu}}$$

$$\Delta y_{\rm avg} \equiv \langle y_{\rm exp} \rangle - \langle y_{\rm QCD} \rangle = \langle E_{\gamma} \rangle / E_{\nu}$$

So, photon takes energy from the charged lepton to the shower, increasing <y>





(Plestid, <u>BZ</u>, 2303.08984)

Correction increases with energy, up to 25%(!)

Final state radiation impacts high-energy (~100 GeV—100 PeV) neutrino observations

Measurements based on inelasticity measurement

- 1. Neutrino-antineutrino flux ratio (5% shift)
- 2. Neutrino mixing parameters
 - Inel. dist. helps to separate nu/nubar
- 3. Charm production from nu interactions
 - CCDIS /w charm production has higher inelasticity

Other measurements

- 1. Throughgoing muons
 - Without FSR, underestimate parent Eν (~5%)
- 2. Double bang signature from tau neutrinos
 - Inference of the parent neutrino energy
 - Reduce the detectability

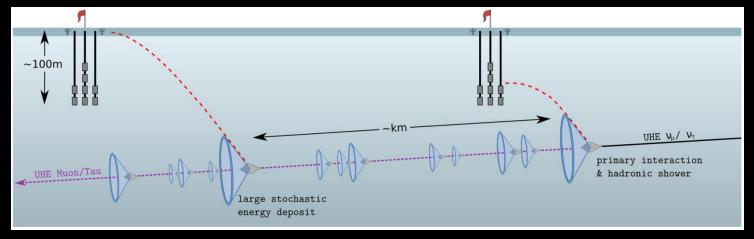
FSR impacts UHE nu observations: in-ice radio detectors (e.g., ANITA, PUEO)

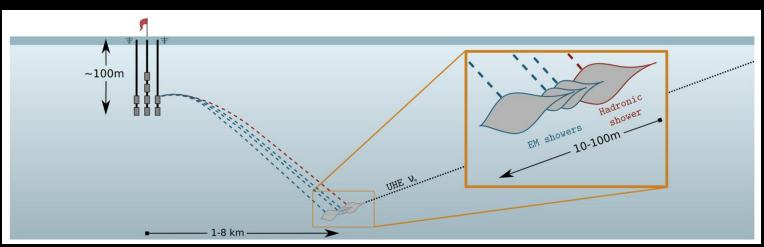
For CCDIS, FSR enhances the overall detectable (shower) energy by as much as ≃20% and lowers the energy thresholds.

ντ CC, big, up to $\simeq 20\%$

νμ CC, mild

ve CC, negligible



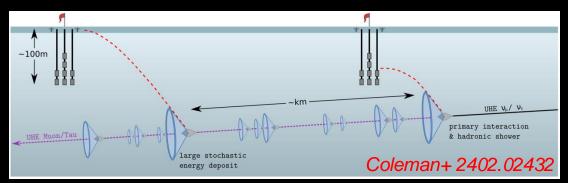


2402.02432 Coleman et al.

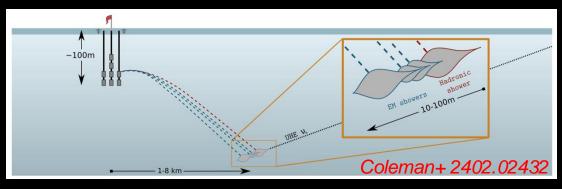
Final state radiation impacts ultrahigh-energy (>~100 PeV) neutrino observations

In-ice radio detectors (e.g., ANITA, PUEO)

- 1) If charged lepton barely detectable, FSR enhances detectable (shower) energy by ~20%.
- 2) A way to measure $\nu\mu/\nu\tau$, FSR reduces the detectability (~5%)



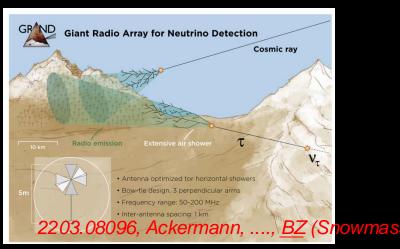
3) A way to measure ve (using LPM effect);



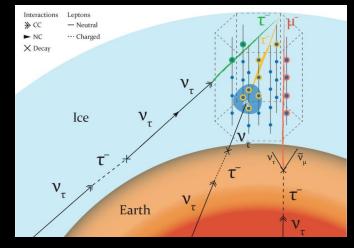
Bkgd rate: 0% (w/o FSR) VS ~30% (w/ FSR)

Air shower detectors (main for $v\tau$, e.g., POEMMA)

1) Earth emergent τ ; w/o FSR underestimates parent Ev by ~5%



2) nutau regeneration, 5%*N



FSR impacts on the neutrino flux and spectrum measurement

Flux normalization:

Any bias on the total detectable energy due to FSR in the previous slides will be amplified when measuring the neutrino flux normalization due to the steeply falling spectrum

$$(1 - \delta_E)^{-\Gamma} \simeq 1 - \Gamma^* \delta_E$$

For example, $\Gamma=3$, $\delta_E=5\%$, the bias is 15% $\Gamma=3$, $\delta_E=20\%$ (UHE $v\tau$ CCDIS), the bias is 60%

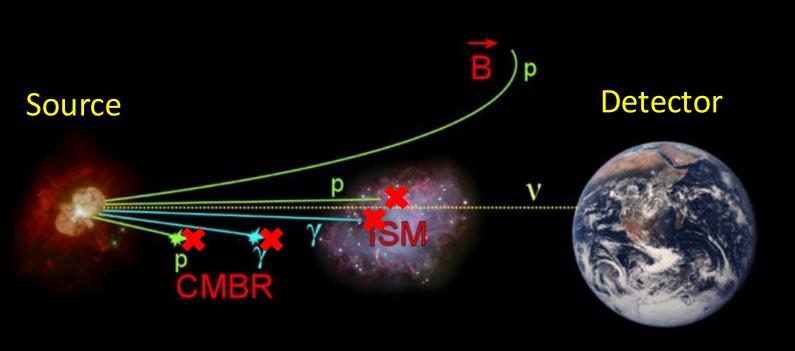
Spectral shape:

FSR's effect is energy dependent, so it affects the spectrum shape measurement. E.g., for IceCube measurements using muon tracks, $\Gamma = -2.37 \pm 0.09 \rightarrow -2.28$ after including FSR. Comparable to the current measurement uncertainty.

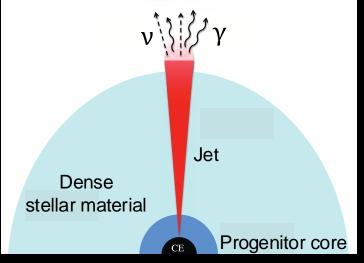
Thanks for your attention!

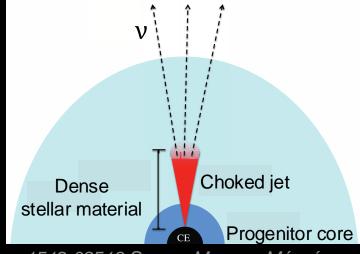
Why do we study high-energy neutrinos: astrophysics

Cosmic ray sources



Dense environment



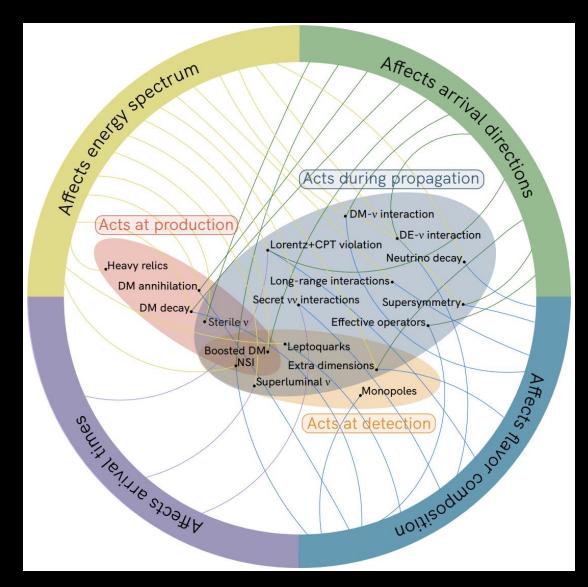


1512.08513 Senno, Murase, Mészáros 2210.03088 Chang, <u>BZ</u>, Murase, Kamionkowski

Why do we study high-energy neutrinos: BSM

Why HE neutrinos special for BSM:

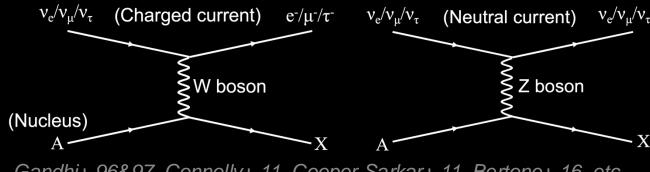
- High energy, inaccessible by lab ν experiments
- **Known direction**
- Travel cosmic distance, small effects accumulates to big effects
- Extremely high column density (through Earth)



HE/UHE neutrino interaction studies so far, not enough

Deep inelastic scattering (DIS) dominates

(~1% precision)



Gandhi+ 96&97, Connolly+ 11, Cooper-Sarkar+ 11, Bertone+ 16, etc.

Most recent: Xie, et al. 2303.13607

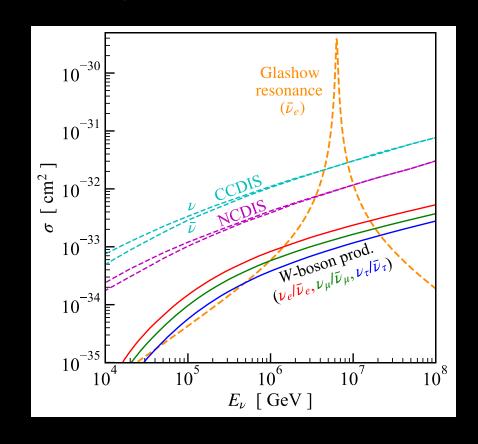
W-boson production is subdominant

Seckel 1997; Alikhanov 2015; <u>BZ</u>, Beacom 1910.08090, 1910.10720

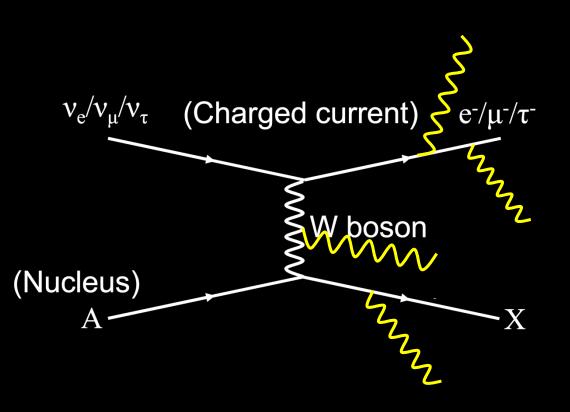
Glashow resonance important for $\bar{\nu}_e$

Glashow 1960, IceCube 2021

Cross sections



Photons from other parts of the diagram: not important



Photon from W boson: suppressed by W mass

Photon from quarks:

- 1) hard to distinguish from the hadronic cascade
- 2) Eγ small as quark energy << lepton energy

Multi-photon emission: higher order, small

A rough estimate using Sudakov form factor

Collinear log Soft log

$$\left| F_S(s, E_{\min}) \sim \exp\left[-\frac{\alpha}{2\pi} \log\left(\frac{s}{m_\ell^2}\right) \log\left(\frac{E_\ell^2}{E_{\min}^2}\right) \right] \right|$$

which gives the probability to *not* radiate any photons above E_{\min} in a collision with center-of-mass energy \sqrt{s} and final-state charged-lepton energy E_{ℓ} . Taking ℓ as the muon (μ) , $E_{\min} \simeq \frac{1}{10} E_{\mu}$, and $s \simeq 2 E_{\nu} m_N$ (m_N) is the nucleon mass) with $E_{\nu} = 10$ TeV, we find $F_S \sim 0.9$. This implies that roughly 10% of all events will contain some prompt real and energetic photon radiation.

Calculation

DIS cross section

$$\frac{\mathrm{d}^2 \sigma_{\nu, \overline{\nu}}^{(0)}}{\mathrm{d}x \mathrm{d}y} = \frac{G_F M E_{\nu}}{\pi (1 + Q^2 / M_W^2)^2} \times \left[y^2 F_1 + (1 - y) F_2 \pm xy (1 - y/2) F_3 \right]$$

from Xie et al. 2303.13607, CTEQ collaboration

Collinear log

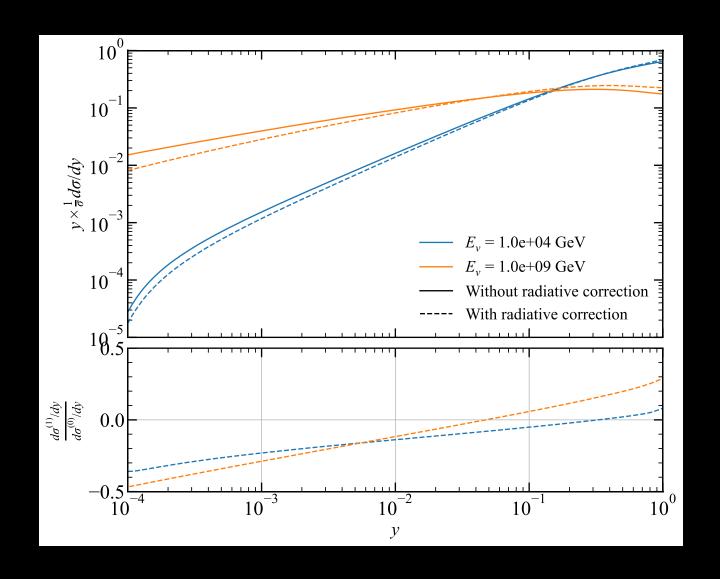
$$P_{\ell \to \ell \gamma}(z) = \frac{\alpha}{2\pi} \log \left(\frac{s}{m_{\ell}^2}\right) \left[\frac{(1+z^2)}{[1-z]_+} + \frac{3}{2}\delta(1-z)\right],\tag{6}$$

$$\frac{\mathrm{d}\sigma^{(1)}}{\mathrm{d}E_{\ell}} = \frac{\alpha}{2\pi} \int \mathrm{d}y \int \mathrm{d}z \, \frac{\mathrm{d}\sigma^{(0)}}{\mathrm{d}y} \delta(E_{\ell} - (1 - y)zE_{\nu})
\times \log\left(\frac{s}{m_{\ell}^{2}}\right) \left[\frac{1 + z^{2}}{[1 - z]_{+}} + \frac{3}{2}\delta(1 - z)\right].$$
(7)

(Plestid, <u>BZ</u>, 2303.08984)

Illustration of FSR impacts on DIS differential xsec

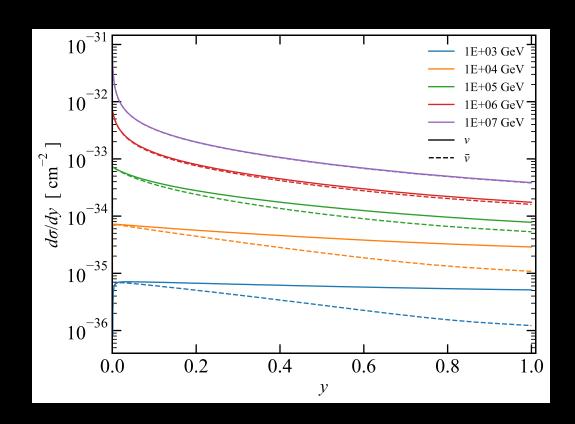
$$y_{\rm QCD} \equiv \frac{E_X}{E_\nu} = \frac{E_\nu - E_\nu}{E_\nu}$$

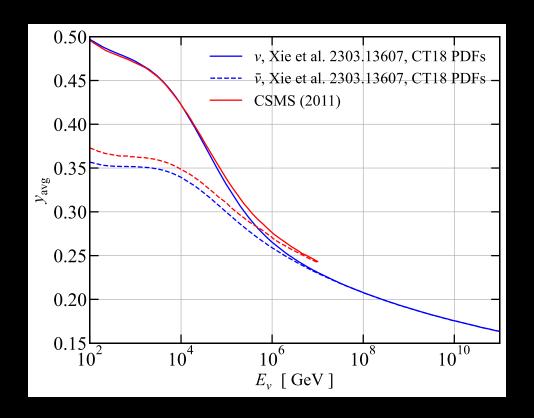


FSR impacts on the inelasticity

Theoretical definition:

$$y_{\rm QCD} \equiv \frac{E_X}{E_\nu} = \frac{E_\nu - E_\ell}{E_\nu}$$





FSR impacts HE nu observation: nu mixing parameters & charm production

Neutrino mixing

Inelasticity measurements help to separate nu and nubar, which helps with measuring neutrino mass hierarchy and CP violation. The sensitivity can be increased by $\approx 30\%$. (1303.0758, 1312.0457, 2402.13308)

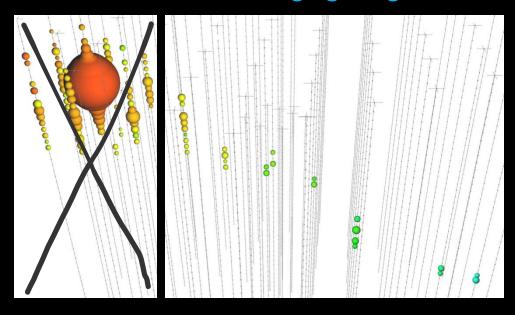
And FSR will affect the measurements

Charm production

Neutrino DIS with charm production has a larger inelasticity than those without...

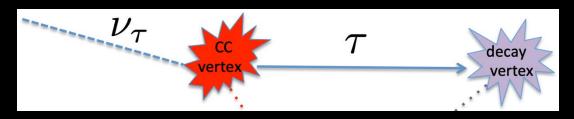
FSR impacts HE nu observation: throughgoing muons & $\nu_{ au}$ double bang

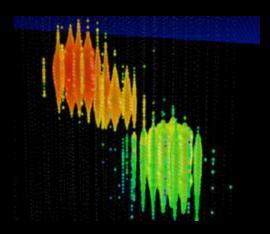
Throughgoing muon



Not including FSR underestimates the parent neutrino energy

ν_{τ} induced double bang

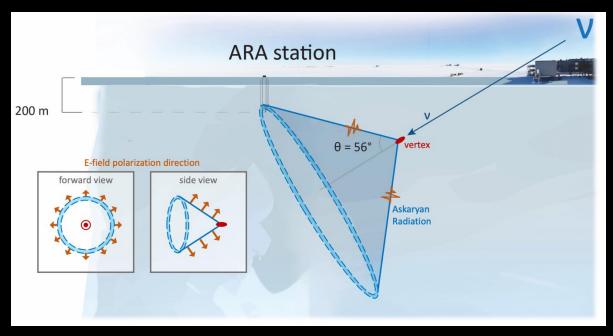




FSR 1) distort the energy balance the two bangs 2) reduce the detectability of the double bang signature.

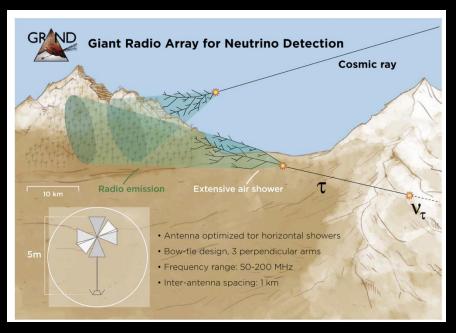
HE nu observation: two basic kinds of detectors

In-ice radio detectors
(all flavors; hard to distinguish flavors)



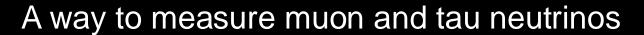
1912.00987 ARA collaboration

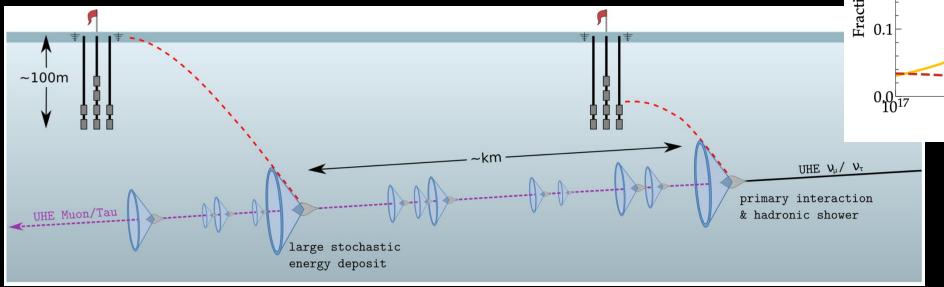
Air shower detectors (main for $v\tau$)

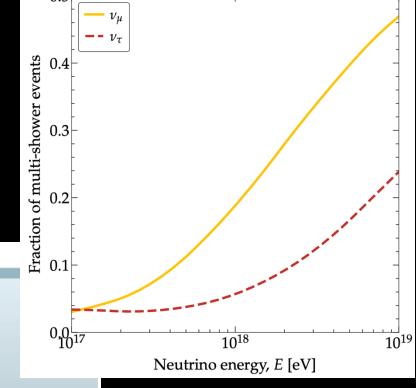


2203.08096, Ackermann,, <u>BZ</u> (Snowmass WP)

FSR impacts UHE nu observation: in-ice radio detectors



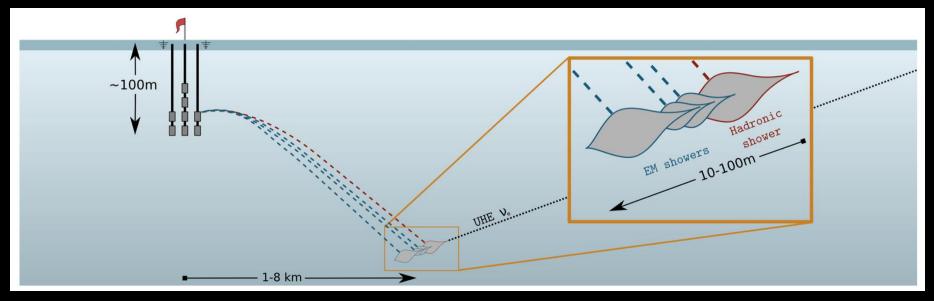




2402.02432 Coleman et al

FSR impacts UHE nu observation: in-ice radio detectors

A way to measure electron neutrinos (using LPM effect)

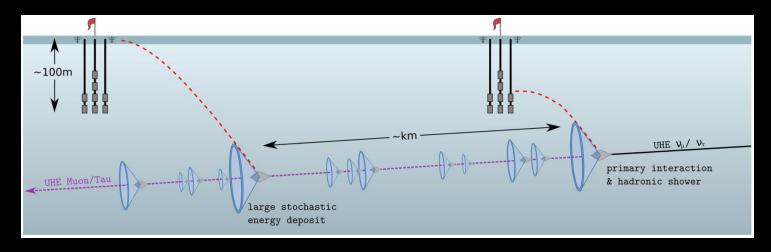


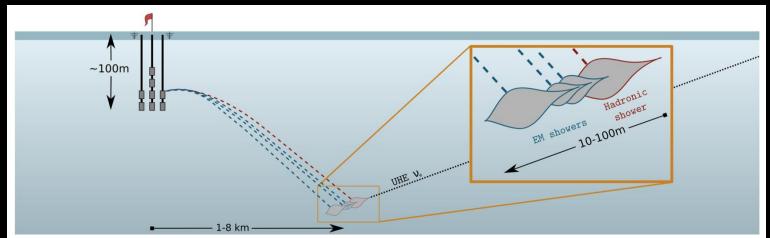
Background could be from muon/tau neutrino CC interactions. Without FSR, the paper estimates that bkgd rate is negligible With FSR, we estimate that bkgd rate is ~30% of signal rate

2402.02432 Coleman et al

FSR impacts UHE nu observation: in-ice radio detectors

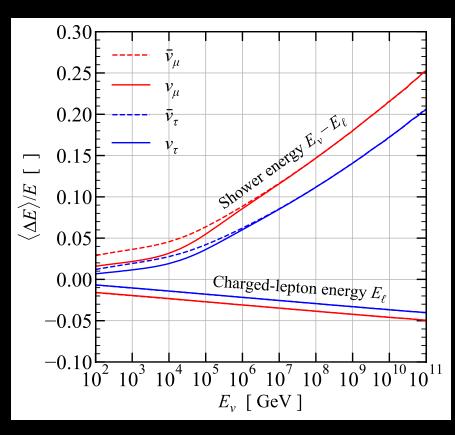
If the charged from CC interaction barely deposit energies to the antenna, then FSR enhances the detectable (shower) energy by as much as 25%!



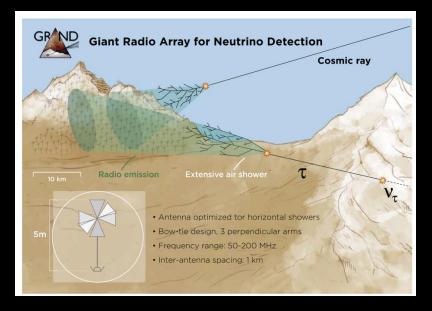


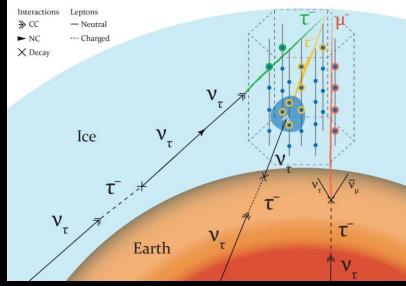
2402.02432 Coleman et al

FSR impacts UHE nu observation: air shower detectors for ντ

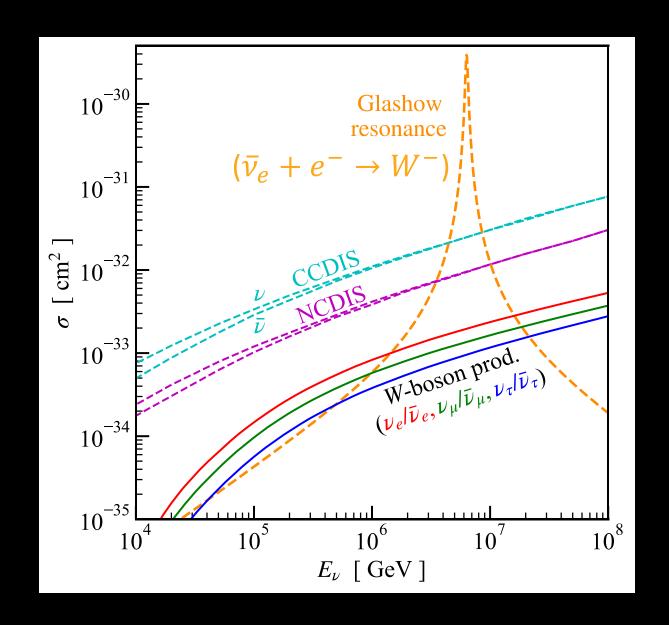


(Plestid, <u>BZ</u>, 2303.08984)





2203.08096, Ackermann, ..., <u>BZ</u> (Snowmass WP)



FSR impacts HE nu detection in collider/accelerator neutrinos

Example: measuring parton distribution function (PDF) using data of FASERv (running) and future FASERv2

FASERv (running) will have ~2×10⁴ neutrino CCDIS events FASERv2 (proposed) will have ~10⁶. Enough data to perform PDF(x, Q²) measurements

Without FSR:
$$x_{(0)} = \frac{Q_{(0)}^2}{2m_N E_X}; \quad Q_{(0)}^2 = 4E_{\nu} E_{\ell} \sin^2\left(\frac{\theta_{\ell}}{2}\right)$$

With FSR:
$$\frac{\Delta Q^2}{Q_{(0)}^2} \simeq -\frac{E_\gamma}{E_\ell}$$
 A few percent but large statistics

$$\frac{\Delta x}{x_{(0)}} \simeq -\frac{E_{\gamma}}{E_X} - \frac{E_{\gamma}}{E_{\ell}}$$
 ~10%